

Exam statistical mechanics (19/12/2023)

① Dieterici equation of state

$$P = \frac{Nk_B T}{V - Nb} e^{-aN/k_B T V}$$

a) $\frac{N}{V} \equiv n$

$$P = \frac{nk_B T}{1 - nb} e^{-an/k_B T}$$

$$\approx nk_B T \left(1 - \frac{an}{k_B T}\right) (1 + nb) + \mathcal{O}(n^3)$$

$$= nk_B T (1 + nb) - an^2 \quad (1)$$

In first order of n one retains

$$P = nk_B T$$

which is the ideal gas law

b) $P_{\text{von}} = \frac{Nk_B T}{V - Nb} - \frac{aN^2}{V^2}$

$$\approx nk_B T (1 + nb) - an^2$$

which matches eq. (1).

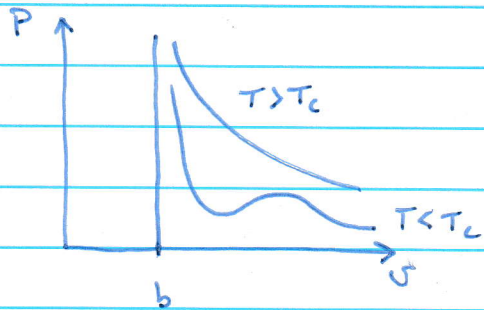
c) System becomes thermodynamically unstable when

$$\left. \frac{\partial P}{\partial V} \right|_T > 0.$$

$$\Leftrightarrow \left. \frac{\partial P}{\partial V} \right|_T = \frac{-Nk_B T}{(V - Nb)^2} \cdot \exp(\dots) + \frac{Nk_B T}{V - Nb} \exp(\dots) \cdot \frac{aN}{k_B T V^2} > 0$$

$$\Leftrightarrow \frac{aN(V - Nb)}{k_B V^2} > T \quad \text{we identify } T_c = \frac{aN(V - Nb)}{k_B V^2}$$

d)



② Quantum partition function

a)

$$Z = \sum_{n=0}^{+\infty} (n+1) e^{-\beta(E_0 + n\gamma)}$$

$$= e^{-\beta E_0} \sum_{n=0}^{+\infty} (n+1) x^n \quad x \equiv e^{-\beta\gamma}$$

$$= e^{-\beta E_0} \cdot \frac{d}{dx} \left(\sum_{n=0}^{+\infty} x^{n+1} \right)$$

$$= e^{-\beta E_0} \cdot \frac{d}{dx} \left(x \sum_{n=0}^{+\infty} x^n \right)$$

$$= e^{-\beta E_0} \frac{d}{dx} \left(\frac{x}{1-x} \right)$$

$$= e^{-\beta E_0} \left(\frac{(1-x) + x}{(1-x)^2} \right)$$

$$= \frac{e^{-\beta E_0}}{(1 - e^{-\beta\gamma})^2}$$

$$\langle E \rangle = -\frac{\partial}{\partial \beta} \log(Z) = E_0 + 2 \frac{\partial}{\partial \beta} \log(1 - e^{-\beta\gamma})$$

$$= E_0 + \frac{2\gamma e^{-\beta\gamma}}{1 - e^{-\beta\gamma}}$$

$$= E_0 + \frac{2\gamma}{e^{\beta\gamma} - 1}$$

low temperature: ($\beta \gg 1$)

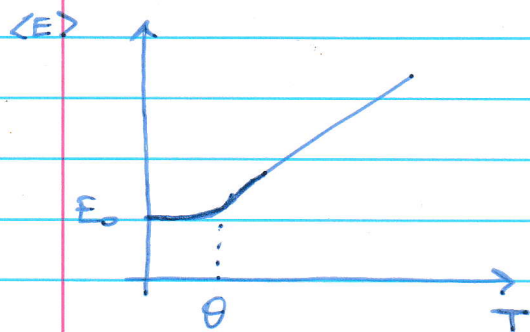
$$\langle E \rangle \approx E_0 + 2\gamma e^{-\beta\gamma}$$

High temperature: ($\beta \ll 1$)

$$\langle E \rangle = E_0 + \frac{2\gamma}{\gamma} \cdot k_B T$$

$$= E_0 + 2k_B T$$

(Insight a 2D harmonic oscillator has exactly this described energy spectrum, hence the $2k_B T$.)



$$b) k_B \theta = \Delta E = \gamma$$

$$\Rightarrow \theta = \frac{\gamma}{k_B}$$

③ Identical particles wave function

a)

$$\Psi(\vec{q}_1, \vec{q}_2, \vec{q}_3) = \frac{1}{\sqrt{3}} \left\{ \phi_1(\vec{q}_1) \phi_2(\vec{q}_2) \phi_2(\vec{q}_3) + \phi_1(\vec{q}_1) \phi_2(\vec{q}_2) \phi_1(\vec{q}_3) + \phi_2(\vec{q}_1) \phi_1(\vec{q}_2) \phi_1(\vec{q}_3) \right\}$$

b)

$$\begin{aligned} \log \Xi_{\text{Bosons}} &= \sum_{\gamma} \log \left(\sum_{n_{\gamma}=0}^{+\infty} e^{-\beta(\epsilon_{\gamma}-\mu)n_{\gamma}} \right) \\ &= - \sum_{\gamma} \log (1 - e^{-\beta(\epsilon_{\gamma}-\mu)}) \\ &= - \log(1 - e^{+\beta\mu}) - \log(1 - e^{-\beta(\epsilon-\mu)}) \end{aligned}$$

c) $\langle n_{\gamma} \rangle = -\frac{1}{\beta} \frac{\partial}{\partial \epsilon_{\gamma}} \log \Xi$

$$\begin{aligned} &= -\frac{1}{\beta} \frac{\partial}{\partial \epsilon_{\gamma}} (-\log(1 - e^{-\beta(\epsilon_{\gamma}-\mu)})) \\ &= \frac{1}{\beta} \frac{\beta e^{-\beta(\epsilon_{\gamma}-\mu)}}{1 - e^{-\beta(\epsilon_{\gamma}-\mu)}} = \frac{1}{e^{\beta(\epsilon_{\gamma}-\mu)} - 1} \end{aligned}$$

$$\langle n_0 \rangle = \frac{1}{e^{-\beta\mu} - 1} ; \quad \langle n_1 \rangle = \frac{1}{e^{\beta(\epsilon-\mu)} - 1}$$

$$\frac{\langle n_0 \rangle}{\langle n_1 \rangle} = \frac{e^{\beta(\epsilon-\mu)} - 1}{e^{-\beta\mu} - 1} = 2$$

$$\Rightarrow e^{\beta\epsilon} x - 1 = 2x - 2, \quad \text{where } x \equiv e^{-\beta\mu}$$

$$\Rightarrow x = \frac{1}{2 - e^{\beta\epsilon}}$$

$$\Rightarrow \mu(T) = +k_B T \log(2 - e^{\beta\epsilon})$$

$$2 - e^{\beta \varepsilon} > 0$$
$$\Rightarrow \frac{\varepsilon}{k_B \log(2)} < T$$

④ Blackbody spectrum

$$\langle E \rangle = \int_0^{+\infty} \frac{g(\omega) \hbar \omega}{e^{\beta \hbar \omega} + 1} d\omega$$

$$= AV \hbar \int_0^{+\infty} \frac{\omega^d d\omega}{e^{\beta \hbar \omega} + 1}$$

$$x = \beta \hbar \omega \rightarrow \frac{dx = d\omega}{\beta \hbar}$$

$$= \frac{AV}{\beta(\beta \hbar)^d} \int_0^{+\infty} \frac{x^d dx}{e^x + 1} \sim T^{d+1}$$

⑤ Fermions

$$\log \Xi_{\text{Fermions}} = \sum_{\gamma} \log (1 + e^{-\beta(\epsilon_{\gamma} - \mu)})$$

$$\langle n_{\gamma} \rangle = \frac{0 \cdot 1 + 1 \cdot e^{-\beta(\epsilon_{\gamma} - \mu)}}{1 + e^{-\beta(\epsilon_{\gamma} - \mu)}} = \frac{1}{e^{\beta(\epsilon_{\gamma} - \mu)} + 1}$$

$$\langle n_{\gamma}^2 \rangle = \frac{0^2 \cdot 1 + 1^2 \cdot e^{-\beta(\epsilon_{\gamma} - \mu)}}{1 + e^{-\beta(\epsilon_{\gamma} - \mu)}} = \frac{1}{e^{\beta(\epsilon_{\gamma} - \mu)} + 1}$$

$$\Delta n_{\gamma}^2 = \langle n_{\gamma}^2 \rangle - \langle n_{\gamma} \rangle^2 \quad (\text{or use derivative trick})$$

$$= \frac{1}{e^{\beta(\epsilon_{\gamma} - \mu)} + 1} - \frac{1}{(e^{\beta(\epsilon_{\gamma} - \mu)} + 1)^2} = \frac{1}{x+1} - \frac{1}{(x+1)^2}$$

$$\text{Define } x \equiv e^{\beta(\epsilon_{\gamma} - \mu)}$$

$$\begin{aligned} \left. \frac{d \Delta n_{\gamma}^2}{dx} \right|_{x=x_{\max}} &= 0 = -\frac{1}{(1+x)^2} + 2 \frac{1}{(x+1)^3} \\ &= \frac{2 - (1+x)}{(1+x)^3} \end{aligned}$$

$$\begin{aligned} \Rightarrow x_{\max} &= 1 \Leftrightarrow e^{\beta(\epsilon_{\gamma} - \mu)} = 1 \\ &\Leftrightarrow \epsilon_{\gamma} = \mu \end{aligned}$$

$$\Delta n_{\gamma \max}^2 = \frac{1}{2} - \frac{1}{4} = \frac{1}{4}$$